

# Kerr-lens mode-locked ytterbium-activated orthoaluminate laser

ZHANG-LANG LIN,<sup>1</sup> WEN-ZE XUE,<sup>1</sup> HUANG-JUN ZENG,<sup>1</sup> GE ZHANG,<sup>1</sup>  
PEIXIONG ZHANG,<sup>2</sup> ZHENQIANG CHEN,<sup>2</sup> ZHEN LI,<sup>2</sup> VALENTIN PETROV,<sup>3</sup>  
PAVEL LOIKO,<sup>4</sup> XAVIER MATEOS,<sup>5</sup> HAIFENG LIN,<sup>6</sup> LI WANG<sup>3</sup> AND  
WEIDONG CHEN<sup>1,3,\*</sup>

<sup>1</sup>Fujian Institute of Research on the Structure of Matter, Chinese Academy of Sciences, Fuzhou, 350002 Fujian, China

<sup>2</sup>Department of Optoelectronic Engineering, Jinan University, 510632 Guangzhou, China

<sup>3</sup>Max Born Institute for Nonlinear Optics and Short Pulse Spectroscopy, Max-Born-Str. 2a, 12489 Berlin, Germany

<sup>4</sup>Centre de Recherche sur les Ions, les Matériaux et la Photonique (CIMAP), UMR 6252 CEA-CNRS-ENSICAEN, Université de Caen, 6 Boulevard Maréchal Juin, 14050 Caen Cedex 4, France

<sup>5</sup>Universitat Rovira i Virgili (URV), Física i Cristal·lografia de Materials (FICMA), 43007 Tarragona, Spain

<sup>6</sup>College of Physics and Optoelectronic Engineering, Shenzhen University, 518118 Shenzhen, China

\*Corresponding author: [chenweidong@fjirsm.ac.cn](mailto:chenweidong@fjirsm.ac.cn)

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**We report on the first Kerr-lens mode-locked laser based on an Yb<sup>3+</sup>-doped perovskite-type orthoaluminate crystal exploiting two different principal light polarizations. The Yb:(Y,Gd)AlO<sub>3</sub> laser delivers soliton pulses as short as 32 fs at 1067 nm with an average output power of 328 mW, a pulse repetition rate of ~84.6 MHz for *E* || *a* polarization. For the orthogonal *E* || *b* polarization, 33-fs pulses are generated at 1057 nm with an average output power of 305 mW. Power scaling to a maximum average output power reaching 2.07 W is achieved at the expense of longer pulse duration (72 fs for *E* || *b*), corresponding to an optical efficiency of 43.9% and a peak power of 303 kW. © 2022 Optical Society of America**

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Mode-locked (ML) solid-state Yb-lasers emitting few-optical-cycle pulses in the spectral range of ~1 μm represent a hot topic in the field of ultrafast optics. Compared to ML Ti:sapphire lasers operating near 800 nm, they feature the following advantages: i) extremely low quantum defect resulting in very high optical efficiency and reduced thermal issues; ii) power-scalable operation using cost-effective and commercially available high-power fiber-coupled InGaAs laser diodes as pump sources; and iii) compact configurations suitable for robust integration [1].

Recently, efforts towards exploiting the broadband gain spectra of Yb<sup>3+</sup>-doped disordered crystals for few-optical-cycle pulse generation resulted in entering the sub-20 fs time domain. Y. Wang *et al.* passed this mark for the first time in 2021 using Yb:CaGdAlO<sub>4</sub> aluminate as a gain medium and generating 17.8 fs pulses at 1118 nm with an average output power of only 26 mW (optical efficiency  $\eta_{\text{opt}} = 3.3\%$ ) via Kerr-lens mode-locking (KLM) [2]. Subsequently, 17 fs pulses at 1080 nm were emitted from a similar

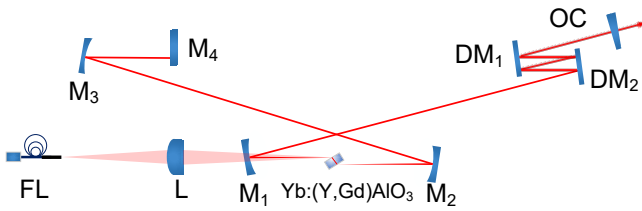
KLM laser based on an isostructural aluminate crystal, Yb:CaYAlO<sub>4</sub>, albeit with a lower average output power of 18 mW and  $\eta_{\text{opt}}$  of only 0.47% [3]. These results indicate quite clearly the existing trade-off between ultimate pulse shortening and output power scalability. To overcome this limitation, a pump scheme allowing one to reduce the quantum defect was employed in a KLM Yb:CaGdAlO<sub>4</sub> laser delivering soliton pulses as short as 22 fs at 1040.7 nm with an average output power of 729 mW and a reasonably high  $\eta_{\text{opt}} = 25\%$  [4]. Ordered host crystals normally possess higher thermal conductivity and are more suited for high output powers/optical efficiencies but classical crystals such as YAG do not offer such broad spectral gain. For instance, a diode-pumped KLM Yb:YAG laser generating 35 fs pulses at 1060 nm with an average output power of 107 mW was demonstrated in 2011 [5]. A shorter pulse duration of 27 fs at 1028 nm for an average output power of 3.3 W was obtained from a KLM thin-disk Yb:YAG laser albeit at an optical efficiency of only 1.05% [6].

Ytterbium (Yb<sup>3+</sup>)-doped yttrium orthoaluminate, Yb:YAlO<sub>3</sub> (shortly Yb:YAP), is a structurally ordered laser crystal belonging to the orthorhombic class (sp. gr. *Pnma*, perovskite-type structure). It is very promising for the development of high-power 1 μm lasers due to a combination of excellent thermo-mechanical and spectroscopic properties [7-10]. In particular, it possesses relatively high thermal conductivity,  $\kappa_a = 7.1$ ,  $\kappa_b = 8.3$  and  $\kappa_c = 7.6$  Wm<sup>-1</sup>K<sup>-1</sup> (values along the crystallographic axes for 5 at.% Yb<sup>3+</sup> doping) [11], which is beneficial for power scalable operation under pumping with commercial high-power InGaAs laser diodes emitting at ~0.98 μm. YAP is optically biaxial and its intrinsic birefringence together with the polarization anisotropy of the Yb<sup>3+</sup> ion emission properties leads to a linearly polarized laser emission, suppressing thermally induced depolarization losses inherent to cubic laser crystals such as Yb:YAG. The first investigation on its passive mode-locking was reported in 2008: using a Semiconductor Saturable Absorber Mirror (SESAM), 225 fs pulses were generated at 1041 nm with an

average output power of 0.8 W [12]. Subsequently, the same group reported on power scaling of the SESAM ML Yb:YAlO<sub>3</sub> laser using a novel off-axis pumping scheme yielding 140 fs pulses at 1009.7 nm with an average output power of 4 W [7]. However, no sub-100 fs ML Yb:YAP laser operation has been reported, yet.

Very recently, we developed a Yb<sup>3+</sup>-doped yttrium-gadolinium “mixed” orthoaluminate crystal exhibiting compositional disorder, i.e., Yb:Y<sub>1-x</sub>Gd<sub>x</sub>AlO<sub>3</sub> or Yb:(Y,Gd)AlO<sub>3</sub> [13]. The estimated mean thermal conductivity of this material is  $\sim 6.3 \text{ Wm}^{-1}\text{K}^{-1}$  [14, 15], indicating a high potential for power scalable laser operation. By implementing a SESAM for starting and stabilizing the soliton pulse shaping, we achieved 43 fs pulses from a diode-pumped Yb:(Y,Gd)AlO<sub>3</sub> laser [15]. The superior thermo-optical properties and preliminary mode-locking results motivated us to further explore the potential of the Yb:(Y,Gd)AlO<sub>3</sub> crystal for few-optical-cycle generation via soft-aperture KLM.

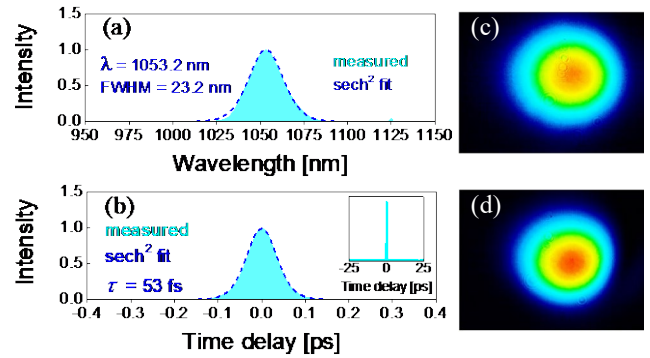
A high-quality Yb<sup>3+</sup>-doped Y<sub>0.85</sub>Gd<sub>0.10</sub>Yb<sub>0.05</sub>AlO<sub>3</sub> crystal with an Yb<sup>3+</sup> doping level of 5.65 at.% was grown by the conventional Czochralski method. The laser operation was investigated using an X-shaped astigmatically compensated standing-wave cavity. The schematic of the laser cavity and pumping geometry is presented in Fig. 1.



**Fig. 1.** Schematic of the Yb:(Y,Gd)AlO<sub>3</sub> laser. FL: fiber laser; L: spherical focusing lens ( $f = 75 \text{ mm}$ ); M<sub>1</sub> – M<sub>3</sub>, dichroic concave mirrors; M<sub>4</sub>: flat rear mirrors; DM<sub>1</sub> and DM<sub>2</sub>: dispersive mirrors; OC: output coupler.

An uncoated 3-mm thick laser crystal was cut from the as-grown bulk, oriented for light propagation along the crystallographic *c*-axis (*c*-cut) with an aperture of 4 (a)  $\times$  4 (b) mm<sup>2</sup> (sp. gr. *P<sub>nm</sub>*). It was mounted in a water-cooled copper holder (coolant temperature: 19°C) and placed at Brewster’s angle between two concave dichroic folding mirrors M<sub>1</sub> and M<sub>2</sub> (radius of curvature: RoC = -100 mm) with the minimum loss condition fulfilled for both the pump and laser wavelengths and polarizations. One arm of the laser cavity was completed by an additional concave mirror M<sub>3</sub> (RoC = -100 mm) and a flat rear mirror M<sub>4</sub>. The other cavity arm was terminated by a pair of flat dispersive mirrors (DMs) and a flat-wedged output coupler (OC). DM<sub>1</sub> and DM<sub>2</sub> were implemented to compensate the intracavity group delay dispersion (GDD) and balance the self-phase modulation (SPM) induced by the Kerr nonlinearity of the laser crystal for soliton pulse shaping. All DMs had a negative GDD of  $-200 \text{ fs}^2$  per bounce. A continuous-wave, spatially single-mode fiber laser emitting a nearly diffraction-limited ( $M^2$  of  $\sim 1.04$ ) linearly polarized beam was used as a pump source. Its emission wavelength was locked at 979 nm with a spectral linewidth (full width at half maximum, FWHM) of  $\sim 0.1 \text{ nm}$ . The pump beam was focused into the laser crystal by a spherical lens ( $f = 75 \text{ mm}$ ) yielding a beam waist radius of  $15.4 \mu\text{m} \times 28.6 \mu\text{m}$  in the sagittal and tangential planes, respectively. By rotating the crystal, we were able to set the laser polarization ( $E \parallel a$  or  $E \parallel b$ ).

The Yb:(Y,Gd)AlO<sub>3</sub> laser was initially characterized for light polarization  $E \parallel b$  corresponding to stronger pump absorption ( $>99\%$  in ML operation) and an OC transmittance ( $T_{OC}$ ) of 4.5%. The KLM regime was investigated at an incident pump power of 3.05 W by applying four bounces (single pass) on the flat DMs giving a round-trip GDD of  $-1600 \text{ fs}^2$ . In order to discriminate the CW and KLM regimes, the laser cavity was aligned close to the edge of the stability region through translating the folding mirror (M<sub>2</sub>) by several hundreds of micrometers away from the pump mirror (M<sub>1</sub>), which resulted in a significant drop of the CW output power. KLM operation could be initiated by a slight knock on the OC or translating the flat rear mirror M<sub>4</sub>. This led to an abrupt increase of the output power from 0.535 to 1 W. The spectrum of the femtosecond pulses was centered at 1053.2 nm with a FWHM of 23.2 nm by assuming a sech<sup>2</sup>-shaped spectral profile, see Fig. 2(a).

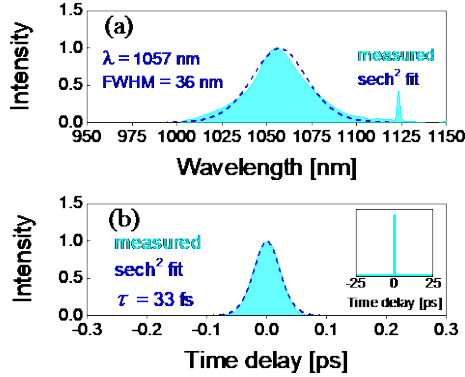


**Fig. 2** KLM Yb:(Y,Gd)AlO<sub>3</sub> laser with  $T_{OC} = 4.5\%$  ( $E \parallel b$ ): (a) Optical spectrum and (b) autocorrelation trace. *Inset* in (b): simultaneously measured long-scale (50 ps) autocorrelation trace. (c,d) Far-field beam profiles: transition from (c) CW to (d) KLM regime.

The recorded background-free intensity autocorrelation trace can be almost perfectly fitted with a sech<sup>2</sup>-shaped temporal profile, giving an estimated pulse duration of 53 fs (FWHM), see Fig. 2(b). The corresponding time-bandwidth product (TBP) of 0.332 was only slightly above the Fourier-transform-limited value (0.315). A long-scale intensity autocorrelation scan of 50 ps confirmed single-pulse steady-state ML, see the inset in Fig. 2(b). The total cavity length was  $\sim 1.8 \text{ m}$ , resulting in a pulse repetition rate of  $\sim 83.2 \text{ MHz}$ . The peak output power reached  $\sim 200 \text{ kW}$ . The calculated peak on-axis intracavity intensity in the laser crystal was  $\sim 1.23 \text{ TW/cm}^2$ . Such a high intensity should lead to a noticeable modification of the spatial mode. This was indeed confirmed by monitoring the far-field beam profiles with a camera placed at  $\sim 1.6 \text{ m}$  away from the OC. The shrinking of the far-field beam diameter from  $2.53 \times 2.23 \text{ mm}^2$  to  $2.17 \times 2.01 \text{ mm}^2$ , as shown in Fig. 2(c) and (d), indicated a strong soft-aperture Kerr-lens effect and self-focusing inside the laser crystal.

The pulse duration could be reduced by applying lower  $T_{OC}$  of 2.5% while maintaining the same total round-trip negative GDD. Notably shorter pulses of 38 fs were obtained at 1056 nm with a  $\sim 31 \text{ nm}$  spectral FWHM. In this case, the corresponding TBP was 0.317, again very close to the Fourier-transform-limit. The average output power was 0.763 W (peak power:  $\sim 212.1 \text{ kW}$ ) at a pulse repetition rate of  $\sim 83.3 \text{ MHz}$ , yielding a peak on-axis laser intensity in the laser crystal of  $\sim 2.37 \text{ TW/cm}^2$ . The shortest pulse duration for laser polarization  $E \parallel b$  was achieved by further reducing the

transmittance of the OC to 1.6%. Soliton pulses as short as 33 fs were generated at 1057 nm with a spectral FWHM of 36 nm, as shown in Fig. 3. The TBP was 0.319, only slightly above the Fourier limit for  $\text{sech}^2$ -shaped pulses. The peak output power of the laser pulses was 97.8 kW corresponding to an estimated peak on-axis laser intensity of  $\sim 1.72 \text{ TW/cm}^2$  in the crystal.



**Fig. 3.** KLM Yb:(Y,Gd)AlO<sub>3</sub> laser with  $T_{oc} = 1.6\%$  ( $E \parallel b$ ): (a) Optical spectrum and (b) autocorrelation trace. *Inset* in (b): long-scale (50 ps) autocorrelation trace.

A maximum output power of 305 mW at  $\sim 83.2 \text{ MHz}$  was obtained for an incident pump power of 3.1 W, corresponding to an optical efficiency of 9.84%. A sharp satellite peak is observed at 1123 nm, which originates from the unmanaged intracavity GDD at the long-wave spectral wing and the non-optimized spectral reflectivity of the cavity mirrors. This phenomenon has been already observed in ML Yb:CaGaAlO<sub>4</sub> lasers [16], and explained by F. Druon *et al.* [17]. It could be improved by properly managing the total intracavity GDD over the full spectral band of the ML laser.

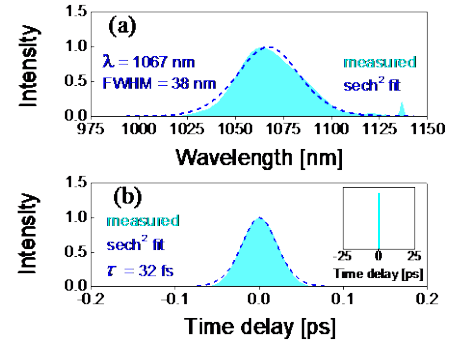
**Table 1. Characteristics<sup>a</sup> of the KLM Yb:(Y,Gd)AlO<sub>3</sub> laser vs Output Coupling for Laser Polarization  $E \parallel b$**

$T_{oc}$ [%]	$P_{pump}$ [W]	$P_{out}$ [W]	$\eta_{opt}$ [%]	$\lambda_c$ [nm]	$\Delta\lambda$ [nm]	$\Delta\tau$ [fs]
1.6	3.1	0.305	9.84	1057	36	33
2.5	3.2	0.763	23.8	1056	30.9	38
4.5	3.05	1	32.8	1053.2	23.2	53
7.5	3.9	1.4	35.9	1053.5	20.7	58
10	4.72	2.07	43.9	1050.1	16	72

<sup>a</sup> $T_{oc}$  – output coupler transmittance,  $P_{pump}$  – incident pump power,  $P_{out}$  – average output power,  $\eta_{opt}$  – optical efficiency refers to incident pump power,  $\lambda_c$  – central laser wavelength,  $\Delta\lambda$  – spectral bandwidth (FWHM),  $\Delta\tau$  – pulse duration (FWHM).

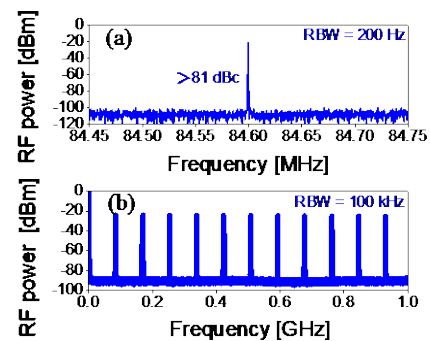
Power scalable operation for the laser polarization  $E \parallel b$  was investigated by increasing the output coupling to  $T_{oc} = 7.5\%$  and 10%. The pump power was adjusted for achieving single-pulse mode-locking with the highest average output power, as well as the highest optical efficiency. Table 1 summarizes the results for various  $T_{oc}$  values. The average output power of the KLM Yb:(Y,Gd)AlO<sub>3</sub> laser was easily scaled by cranking up the pump power at the expense of the pulse duration. The maximum average output power reached 2.07 W at an incident pump power of 4.72 W. This corresponded to a longer pulse duration of 72 fs, still in the sub-100 fs time range, and a relatively high optical efficiency of 43.9%.

The spectrum of the ML laser experienced a slight blue-shift to 1053.5 nm ( $T_{oc} = 7.5\%$ ) and 1050.1 nm ( $T_{oc} = 10\%$ ), as can be expected for a three-level laser system with reabsorption.



**Fig. 4.** KLM Yb:(Y,Gd)AlO<sub>3</sub> laser with  $T_{oc} = 1.6\%$  ( $E \parallel a$ ): (a) Optical spectrum and (b) autocorrelation trace. *Inset* in (b): long-scale (50 ps) autocorrelation trace.

The KLM operation of the Yb:(Y,Gd)AlO<sub>3</sub> laser was also studied for the other available light polarization,  $E \parallel a$ , characterized by lower pump absorption ( $< 75\%$  in ML operation). As expected, the shortest pulses with ultimate mode-locking stability were achieved with the 1.6% OC. A maximum average output power of 328 mW was obtained at an incident pump power of 3.75 W and a pulse repetition rate of 84.6 MHz. This corresponded to an optical efficiency of 8.7%. The pulse characteristics are shown in Fig. 4. The laser spectrum had a FWHM of 38 nm at 1067 nm by assuming a  $\text{sech}^2$ -shaped spectral profile, see Fig. 4(a). The pulse duration was assessed from the recorded background-free intensity autocorrelation trace which is shown in Fig. 4(b). The curve can be almost perfectly fitted with a  $\text{sech}^2$ -shaped temporal profile yielding a deconvolved FWHM of 32 fs. The TBP was thus 0.320, slightly above the Fourier-transform-limit. A longer time window (50-ps) autocorrelation scan confirmed single-pulse mode-locking without post- or pre-pulses, see the inset in Fig. 4(b). The peak output power reached  $\sim 106.6 \text{ kW}$  and the estimated peak on-axis laser intensity in the crystal amounted to  $\sim 1.87 \text{ TW/cm}^2$ .



**Fig. 5.** RF spectra of the KLM Yb:(Y,Gd)AlO<sub>3</sub> laser ( $E \parallel a$ ): (a) Fundamental beat note, RBW = 200 Hz; (b) 1-GHz span with RBW = 100 kHz. RBW: resolution bandwidth.

Figure 5 shows the recorded radio-frequency (RF) spectrum of the steady-state pulse train of the shortest pulses for light

polarization  $E \parallel a$ . The recorded narrow-band fundamental beat note exhibited a very high extinction ratio of 81 dBc above carrier at 84.6 MHz, see Fig. 5(a). This, together with the uniform harmonic beat notes on a 1-GHz frequency span indicates highly stable KLM operation without any Q-switching instabilities or multi-pulsing, see Fig. 5(b).

Power scaling of the KLM laser for light polarization  $E \parallel a$  was also investigated by increasing the output coupling in the range of  $T_{oc} = 2.5\% - 7.5\%$ , cf. Table 2 summarizing the achieved output characteristics. Due to the relatively low pump absorption compared to  $E \parallel b$ , the maximum incident pump power reached 6 W for achieving single-pulse mode locking with higher average output power and optical efficiency. The maximum average output power reached 1.72 W at an incident pump power of 6 W for a pulse duration of 51 fs. This, corresponded to a lower optical conversion of 28.7% compared to  $E \parallel b$ .

**Table 2. Characteristics<sup>a</sup> of the KLM Yb:(Y,Gd)AlO<sub>3</sub> laser vs Output Coupling for Laser Polarization  $E \parallel a$**

$T_{oc}$ [%]	$P_{pump}$ [W]	$P_{out}$ [W]	$\eta_{opt}$ [%]	$\lambda_c$ [nm]	$\Delta\lambda$ [nm]	$\Delta\tau$ [fs]
1.6	3.75	0.328	8.7	1067	38	32
2.5	3.3	0.498	15.1	1065.3	33.3	37
4.5	6	1.24	20.7	1053.5	25.3	50
7.5	6	1.72	28.7	1051.8	22.9	51

<sup>a</sup> $T_{oc}$  – output coupler transmittance,  $P_{pump}$  – incident pump power,  $P_{out}$  – average output power,  $\eta_{opt}$  – optical efficiency refers to incident pump power,  $\lambda_c$  – central laser wavelength,  $\Delta\lambda$  – spectral bandwidth (FWHM),  $\Delta\tau$  – pulse duration (FWHM).

To conclude, “mixed” ytterbium-doped yttrium-gadolinium orthoaluminate crystal, Yb:(Y,Gd)AlO<sub>3</sub>, with compositional disorder, represents an attractive laser material for generation of ultrashort (sub-100 fs) pulses with moderate to high average output powers (up to the watt-level) from Kerr-lens mode-locked solid-state lasers. Owing to its good thermal properties, as well as intense and broad emission bands for polarized light, two regimes of KLM operation were successfully exploited: i) targeting the shortest pulse duration, 32 fs pulses were generated at 1067 nm with an average output power of 328 mW at a repetition rate of 84.6 MHz for light polarization  $E \parallel a$  using low output coupling ( $T_{oc} = 1.6\%$ ), ii) targeting the maximum average output power and laser efficiency, the KLM laser delivered 2.07 W at 1050.1 nm at the expense of longer pulse duration (72 fs) corresponding to an optical efficiency of 43.9% for light polarization  $E \parallel b$  employing high output coupling ( $T_{oc} = 10\%$ ). To the best of our knowledge, this is the first demonstration of Kerr-lens mode-locking employing Yb<sup>3+</sup>-doped orthorhombic perovskite-type aluminate crystals. Our results pave the way towards the development of high-power (>10 W) KLM ytterbium orthoaluminate lasers pumped by cost-efficient spatially multi-mode InGaAs laser diodes. Heavily-doped Yb:YAlO<sub>3</sub>-type crystals also look promising for KLM thin-disk lasers at  $\sim 1 \mu\text{m}$

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**Disclosures.** The authors declare no conflicts of interest

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