

Optimizing the Dynamics of a Two-cell DC-DC Buck Converter by Time Delayed Feedback Control[☆]

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Abstract

A study of the dynamical behavior of a two-cell DC-DC buck converter under a digital time delayed feedback control (TDFC) is presented. Various numerical simulations and dynamical aspects of this system are illustrated in the time domain and in the parameter space. Without TDFC, the system may present many undesirable behaviors such as sub-harmonics and chaotic oscillations. TDFC is able to widen the stability range of the system. Optimum values of parameters giving rise to fast response while maintaining stable periodic behavior are given in closed form. However, it is detected that in a certain region of the parameter space, the stabilized periodic orbit may coexist with a chaotic attractor. Boundary between basins of attraction are obtained by means of numerical simulations.

Keywords: Switching converters, Two-cell buck converter, TDFC,

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1 **1. Introduction**

2 Since the middle of the last century, electrical power has become the most
3 appealing power due to its easiness to be used in industrial as well as resi-
4 dential environments [1, 2]. Therefore, the need to develop ways to convert
5 and transmit this kind of power has drawn the interest of several researchers.
6 During the last few decades, power electronics has undergone a boosting de-
7 velopment in many aspects of technology, including power devices, circuit
8 design, control methods and computer-aided analysis [3]. The main focus in
9 power electronics circuit design remains the fulfilment of the functional re-
10 quirements of the application for which the circuit has been designed. Power
11 electronics converters have been widely used for more than three decades.
12 During this period, several unusual strange behaviors have been observed
13 (*e.g.* subharmonic oscillations, bifurcations and chaos) [4, 5] but very often
14 they have been avoided by adjusting component values and parameters or by
15 adding auxiliary security circuits [6, 7]. During the last two decades, tools for
16 analyzing bifurcations and chaos have been well developed [8, 9], therefore,
17 making the investigation of their very peculiar aspect an attractive endeavor
18 [10, 11, 12].

19 Analysis, characterization and modeling of simple switching converters
20 are still actively pursued today [13, 14]. It has been discovered that de-
21 spite their simplicity, power converters are copious in nonlinear dynamics.
22 This property is owed to toggling between a set of linear or nonlinear circuit

23 topologies, under the control of a feedback control system. Therefore, the
24 control method as well as the circuit topology directly influence the dynam-
25 ical behavior of the power converter [15].

26 Simple DC-DC converters, in all their simple types: buck, buck-boost
27 and boost, have been thoroughly studied in recent years and different types
28 of bifurcations have been perceived [9]. More recently, there has been an
29 interest in analyzing the nonlinear behavior of the multi-cell power electronic
30 converters [18, 19]. As a matter of fact, multi-cell power converters have
31 been widely used in real applications since they circumvent shortcomings
32 of ordinary switching devices due to their ability to support high voltages
33 [20]. Recently, a two-cell buck converter has been also proposed for efficient
34 wide-bandwidth envelope tracking in radio frequency power amplifiers [21].
35 It has been shown that chaotic behavior in these systems usually emerges
36 from bad tuning of the feedback gain while trying to eliminate the static
37 error by increasing this gain.

38 With the purpose to avoid the chaotic behavior in nonlinear systems and
39 maintaining stable periodic regime, researchers tried to apply some tech-
40 niques for controlling chaos [22]. The principle of these control methods is
41 based on adding an extra input (preferably low energy signal) to the non-
42 linear system with the aim to modify its dynamics by stabilizing the desired
43 behavior. The control input can be either a certain physical action on the
44 system such as a force in a mechanical system or an external voltage in an
45 electrical circuit or either a variation of some parameter of the controlled
46 system in terms of its state variables trajectory in the state space [23].

47 One of the most efficient chaos control scheme is the Time Delay Feedback

48 Control (TDFC) first introduced by Pyragas [24] and later applied by many
49 other authors. Roughly speaking, the time delayed scheme adds a control
50 input which uses the difference between the feedback signal and its delayed
51 version. In feedback systems, this input is added in the feedback path. If
52 the time delay is selected to be the period of the Unstable Periodic Orbit
53 (UPO), the latter can be stabilized by appropriately choosing the time delay
54 feedback gain. A very important property of this control scheme is that it is a
55 noninvasive control in the sense that the desired behavior of the system is not
56 modified under the action of the TDFC and it is only stabilized. This method
57 has been successfully applied to the elementary DC-DC buck converter [25],
58 [26]. However, its main drawback is its difficulty to be implemented in analog
59 hardware because of the requirement of a continuous memory during one
60 delayed period. A possible solution to this problem is by using a digital
61 TDFC. For this purpose, an accurate discrete time model must be derived
62 for the system in order to be able to study the performances of the system
63 under the digital TDFC [27, 28].

64 The aim of this work is to use a map that describes the nonlinear dy-
65 namical behavior of a two-cell DC-DC buck converter in order to design a
66 TDFC that can achieve stable dynamics within a wide range of the design
67 parameter space. As it will be shown later, the application of the TDFC
68 technique is a powerful strategy for avoiding subharmonic oscillations and
69 chaotic behavior while widening the stability region in the design parameter
70 space.

71 The rest of the paper is outlined as follows. In the next section, we
72 present a complex discrete model of the converter issued from the solutions

73 of the continuous time equations and we present next the simplified model.
 74 Then, we present a series of numerical simulations of the obtained models.
 75 Section three will be devoted to analyzing the effect of adding TDFC. Both
 76 linear stability analysis, based on Jacobian matrix, and nonlinear stability
 77 analysis, based on a Lyapunov function, led to similar results. Coexistence of
 78 attractors are detected from numerical simulations in the same section and
 79 some explanations of this phenomenon are given. Finally, the last section
 80 will include the concluding remarks and the perspectives.

81 **2. Modeling the two-cell DC-DC converter**

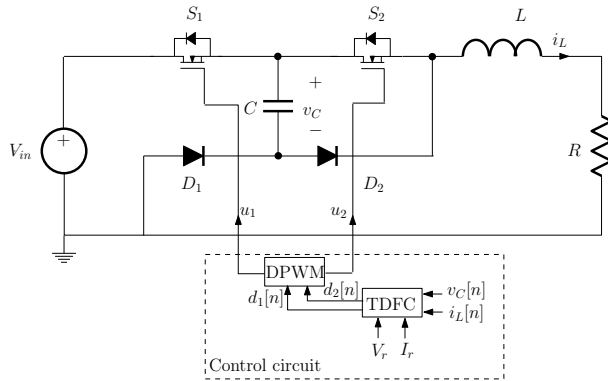


Figure 1: A basic two-cell DC-DC buck converter with an RL load and its control circuit. S_i and D_i are the switch and diode of the i^{th} cell, the capacitor C ensures voltage balance between cells, V_{in} is the supply voltage, v_C and i_L are the circuit states and u_k is the DPWM output signal with duty cycle d_k .

82 The converter that we will deal with in this paper is depicted in Figure
 83 1. It is based on a buck chopper (see [1]) modified in order to allow a higher
 84 input voltage by using two serial switches (transistors and diodes) [18, 19, 20].
 85 The role of the capacitor is to balance the switch voltages. The signals u_1 and

86 u_2 are the outputs of a digital pulse width modulator (DPWM) driven by a
87 feedback control system to be designed in later sections to achieve a constant
88 average value of $v_C = \frac{V_{in}}{2}$ and a constant averaged value of the output current
89 i_L equal to a reference current I_r . The control signals are the duty cycles
90 d_1 and d_2 of the driving signal used for the switches S_1 and S_2 . It is worth
91 noting that in this work, we will define the duty cycles with respect to the
92 OFF state rather than the ON state. Therefore the switch S_k will remain
93 in its OFF state during $d_k T$ when it is driven by a driving signal $u_k(t)$ with
94 a duty cycle d_k . We also notice that the duty cycle d_k cannot exceed unity
95 neither can it go beyond zero, thus when calculated d_k may undergo upper
96 or lower saturation if the obtained value is not in the unit interval.

97 In a typical operating mode, because of the continuity of the inductor
98 current and due to the natural turn off switching of diodes under negative
99 voltages, the switch S_k and the diode D_k are always turned ON and OFF in
100 a complementary manner. Under a fixed frequency PWM control, when the
101 switch S_k is closed, the diode D_k is open and vice versa. As we have two
102 switches S_1 and S_2 , then we can easily define four different operating topolo-
103 gies, (see Table 1). By applying Kirchhoff laws to the two-cell converter, the
104 following switched model is obtained:

$$\frac{d}{dt} \begin{pmatrix} i_L \\ v_C \end{pmatrix} = \begin{pmatrix} -\frac{R}{L} & (u_2 - u_1)\frac{1}{L} \\ (u_1 - u_2)\frac{1}{C} & 0 \end{pmatrix} \begin{pmatrix} i_L \\ v_C \end{pmatrix} + \begin{pmatrix} \frac{V_{in}}{L}u_1 \\ 0 \end{pmatrix} \quad (1)$$

105 where u_k are the driving signals for the switches S_k given by:

$$u_k = \begin{cases} 1 & \text{if } S_k \text{ is closed (ON)} \\ 0 & \text{if } S_k \text{ is open (OFF)} \end{cases} \quad (2)$$

106 Six different operating modes are possible depending on the values of the
 107 duty cycles d_1 and d_2 . The driving signals corresponding to these operating
 108 modes are depicted in Figure 2.

Table 1: Two-cell DC-DC converter different topologies between which toggling occurs

Topologies	state of S_1	state of S_2
Topology 1 (\mathcal{T}_1)	OFF	ON
Topology 2 (\mathcal{T}_2)	ON	ON
Topology 3 (\mathcal{T}_3)	ON	OFF
Topology 4 (\mathcal{T}_4)	OFF	OFF

109 *2.1. Obtaining the discrete time model*

110 Toggling between different topologies occurs according to the values of
 111 the duty cycles d_1 and d_2 at the beginning of the PWM period T . Figure 3
 112 shows the form of steady state the driving signals u_1 and u_2 together with
 113 the their corresponding system response in terms of the state variables i_L
 114 and v_C .

115 As a matter of fact, we can define several different toggling modes, to
 116 each of which we can assign a recurrent system to express the values of the
 117 states at the end of the period in terms of the values at the beginning of the
 118 period. Once at T -periodic steady state, only one recurrent function will
 119 represent the behavior of the converter. However, during the transient phase,
 120 the duty cycles may saturate and may assume any value depending on the
 121 control system and therefore different periods might be governed by various
 122 modes. The two-cell converter can be described by the following switched
 123 mode model:

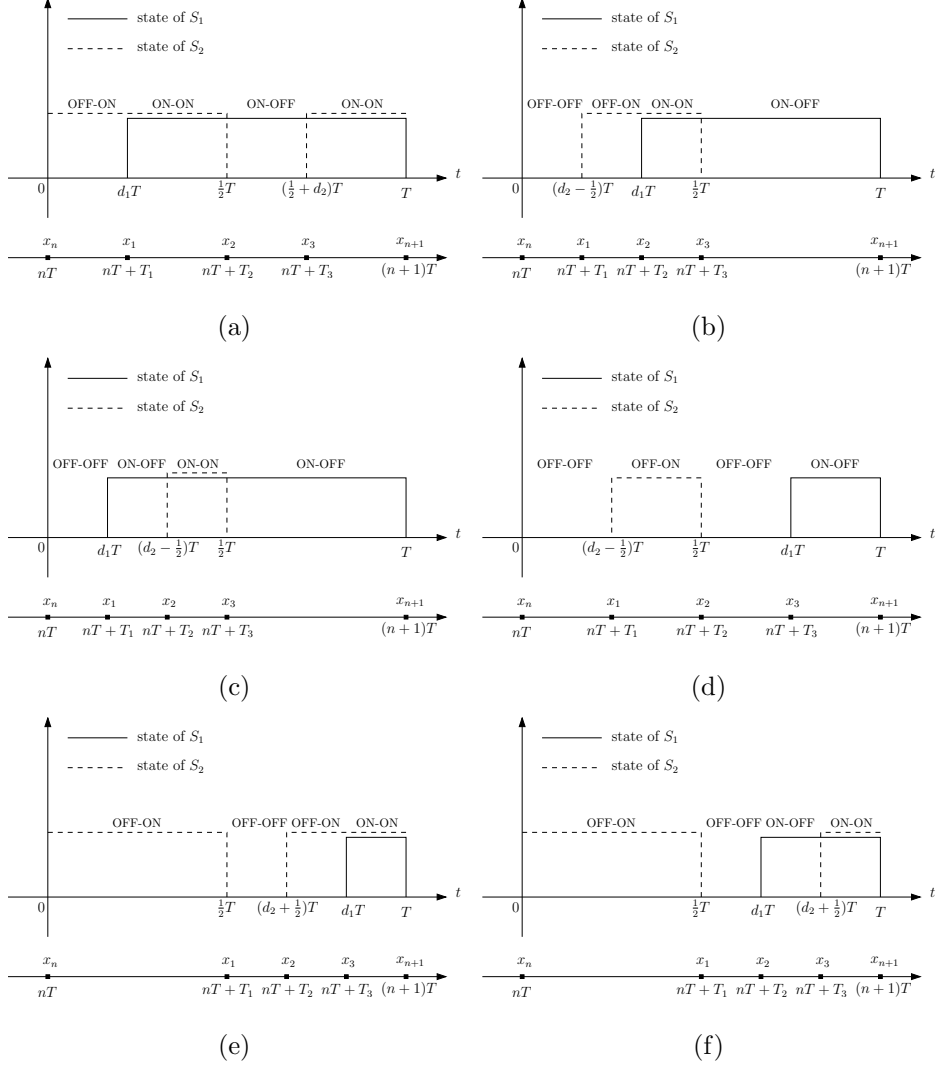


Figure 2: Different possible operating modes for the two-cell DC-DC buck converter. (a) Mode 1: $d_1 < \frac{1}{2}$ and $d_2 < \frac{1}{2}$, (b) Mode 2: $d_1 < \frac{1}{2}$ and $0 < d_2 - \frac{1}{2} < d_1$, (c) Mode 3: $d_1 < \frac{1}{2}$ and $d_1 < d_2 - \frac{1}{2} < \frac{1}{2}$, (d) Mode 4: $d_1 > \frac{1}{2}$ and $d_2 > \frac{1}{2}$, (e) Mode 5: $\frac{1}{2} < d_2 + \frac{1}{2} < d_1$, (f) Mode 6: $d_1 < d_2 + \frac{1}{2} < 1$ and $\frac{1}{2} < d_2 + \frac{1}{2} < d_1$

$$\begin{aligned}
 \dot{x} &= A_1 x + B_1 && \text{during } [t_0, t_1] \\
 \dot{x} &= A_2 x + B_2 && \text{during } [t_1, t_2] \\
 \dot{x} &= A_3 x + B_3 && \text{during } [t_2, t_3] \\
 \dot{x} &= A_4 x + B_4 && \text{during } [t_3, t_4]
 \end{aligned} \tag{3}$$

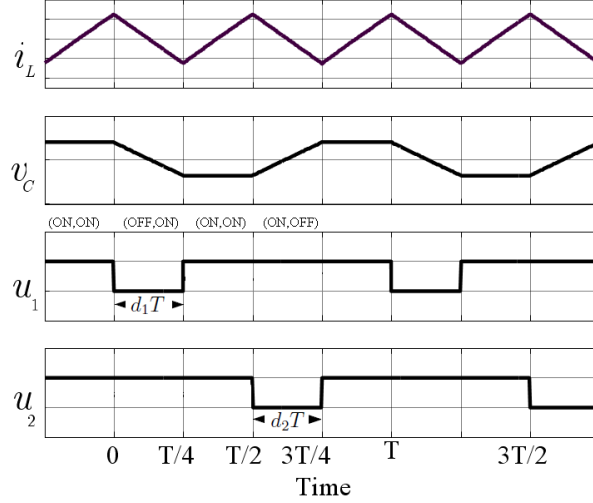


Figure 3: Driving signals u_1 and u_2 and the state variables i_L and v_C of a basic two-cell DC-DC buck converter in steady state periodic behavior.

where A_k and B_k can be easily deduced from the switched model (1) by specifying the operating mode and giving values to the driving signal $u(t)$, t_k ($k = 1 \dots 4$) are the switching instants, being $t_0 = nT$, $t_4 = (n+1)T$ and $(t_1, t_2$ and $t_3)$ decided by the controller (See Fig. 2). Eq (3) is piecewise linear and taking advantage of the linearity of the system during each switching interval, the solution during this interval can be written as follows [16]

$$x(t) = e^{A_k(t-t_k)}x(t_k) + \int_{t_k}^t e^{A_k(t-\tau)}B_k d\tau \quad (4)$$

By stacking-up the solutions, the following model is obtained [18]:

$$x[n+1] = \mathbf{\Phi}(\mathbf{d}[n])x[n] + \mathbf{\Psi}(\mathbf{d}[n]) \quad (5)$$

124 where $x[n] = x(nT)$, $x[n + 1] = x((n + 1)T)$ and

$$\mathbf{d}[n] = (d_1[n], d_2[n])^t \quad (6)$$

$$\Phi(\mathbf{d}[n]) = \phi_4\phi_3\phi_2\phi_1 \quad (7)$$

$$\Psi(\mathbf{d}[n]) = \phi_4\phi_3\phi_2\psi_1 + \phi_4\phi_3\psi_2 + \phi_4\psi_3 + \psi_4 \quad (8)$$

$$\phi_k = e^{A_k\delta t_k} \text{ and } \psi_k = \int_{t_k}^{t_{k+1}} e^{A_k(t-\tau)} B_k d\tau \quad (9)$$

where $\delta t_k = t_{k+1} - t_k$ is the lapse of time between two consecutive switching instants. As it can be noticed, this model is extremely complicated to handle from a theoretical analysis point of view. In order to simplify the nonlinear model of the system while maintaining practical considerations, we will assume some realistic assumptions with the purpose of obtaining an affine expression in the control signals d_1 and d_2 . We will proceed by approximating the entries of the matrix exponential e^{At} by using Taylor's expansion to the first order. This is possible since the time durations between switchings are much smaller than time constant of the system and the state variables can therefore be considered practically piecewise linear. Therefore, it can be written

$$e^{A_k t} \simeq \mathbb{I} + A_k t + \mathcal{O}((A_k t)^2) + \dots \quad (10)$$

With the above consideration, it can be demonstrated that all the above operating modes can be described by the following simplified model [30]

$$x[n + 1] = (\mathbb{I} + AT)x[n] + BT \quad (11)$$

where

$$A = \sum_{k=1}^4 A_k \delta t_k, \quad B = \sum_{k=1}^4 B_k \delta t_k \quad (12)$$

Now by substituting A_k , B_k and δt_k by their expressions, we obtain

$$\begin{pmatrix} i_L[n+1] \\ v_C[n+1] \end{pmatrix} = \begin{pmatrix} 1 - \frac{TR}{L} & (d_1[n] - d_2[n])\frac{T}{L} \\ (d_2[n] - d_1[n])\frac{T}{C} & 1 \end{pmatrix} \begin{pmatrix} i_L[n] \\ v_C[n] \end{pmatrix} \quad (13)$$

$$+ \begin{pmatrix} \frac{V_{in}T}{L}(1 - d_1[n]) \\ 0 \end{pmatrix} \quad (14)$$

For the sake of reducing the number of parameters of the system we will consider some new dimensionless variables, where the current is normalized by the maximum current in the circuit, the voltage is normalized by the maximum voltage and time is normalized by the switching period. If we let $\delta_L = \frac{RT}{L}$, $\delta_C = \frac{T}{RC}$, $x_i[n] = \frac{Ri_L[n]}{V_{in}}$ and $x_v[n] = \frac{v_C[n]}{V_{in}}$, we obtain the following dimensionless model

$$\begin{pmatrix} x_i[n+1] \\ x_v[n+1] \end{pmatrix} = \begin{pmatrix} 1 - \delta_L & (d_1[n] - d_2[n])\delta_L \\ (d_2[n] - d_1[n])\delta_C & 1 \end{pmatrix} \begin{pmatrix} x_i[n] \\ x_v[n] \end{pmatrix} \quad (15)$$

$$+ \begin{pmatrix} \delta_L(1 - d_1[n]) \\ 0 \end{pmatrix}$$

125 In this section we present some numerical simulations with a simple pro-
 126 portional control strategy. Herein, the duty cycles are calculated on the basis
 127 of

$$d_1[n] = k_i(x_i[n] - I_r) + k_v(x_v[n] - V_r) \quad (16)$$

$$d_2[n] = k_i(x_i[n] - I_r) - k_v(x_v[n] - V_r) \quad (17)$$

where k_i and k_v are the proportional gains of the current and voltage errors respectively. We implicitly understand from any expression of the duty cycles d_1 and d_2 that it undergoes a saturation because during one switching period

the OFF time can not be bigger than this period neither negative. Therefore, the applied duty cycle is in fact $\text{sat}(d_k)$ and not d_k , where $\text{sat}(d)$ is the function defined by:

$$\text{sat}(d) = \frac{1}{2} (1 + |d| - |d - 1|) \quad (18)$$

128 and the actual duty cycles become

$$\begin{pmatrix} d_1[n] \\ d_2[n] \end{pmatrix} = \text{sat} \left(\begin{pmatrix} k_i & k_v \\ k_i & -k_v \end{pmatrix} \begin{pmatrix} x_i[n] - I_r \\ x_v[n] - V_r \end{pmatrix} \right) \quad (19)$$

which can be written in a more compact form as $\mathbf{d}[n] = \text{sat}(K(x[n] - X_r))$, where $x[n] = (x_i[n], x_v[n])^t$, $X_r = (I_r, V_r)^t$ and

$$K = \begin{pmatrix} k_v & k_i \\ k_v & -k_i \end{pmatrix} \quad (20)$$

129 Due to the piecewise aspect of the duty cycles, we can easily verify that
 130 the system can mathematically have several fixed points. However, only the
 131 non saturated fixed point has a practical interest. The expression of this
 132 fixed point will be given later.

133 2.2. Model Verification from Numerical simulations

134 In the previous section, we have presented different models of the two-cell
 135 DC-DC buck converter. The first is the exact model (1) which is not appropriate
 136 to deal mathematically with it. Indeed, the model that represents
 137 the overall dynamics of the converter is a non-autonomous switched and non
 138 affine in the control signal. Therefore, the only use of this complex model
 139 is to check via numerical simulations whether any suggested control strategy
 140 would be efficient or not on the exact model, which gives a closer idea to

141 the reality. The second model (5) is the exact discrete time model which has
 142 the advantage of carrying out fast numerical simulations but it presents the
 143 same shortcoming of the switched model in terms of mathematical treatment.
 144 The third model (15) which comes from tenable simplification of the exact
 145 discrete time model presents several advantages. First of all, the model is
 146 unique for all operating continuous conduction mode¹. Second, the obtained
 147 model is simple making its analysis easier.

148 2.2.1. Simulating the simple model

The normal periodic running mode of the controlled converter, which is the period-one mode, corresponds to the fixed point of the recurrent relation (15) together with (16)-(17). The fixed point denoted by (x_i^*, x_v^*) is obtained by solving $(x_i[n+1], x_v[n+1]) = (x_i[n], x_v[n])$. This yields to

$$(x_i^*, x_v^*) = \left(\frac{1 + k_i I_r}{1 + k_i}, V_r \right) \quad (21)$$

It is clear that while, under mild conditions, the reference voltage will be attained, the reference current can only be approached from above when k_i is very large. Mathematically written

$$\lim_{k_i \rightarrow \infty} x_i^* = I_r^+ \quad (22)$$

149 An in-depth analysis of conditions on the stability of the closed loop system
 150 will be elucidated in the next section.

151 The chaotic behavior is well delineated using the bifurcation diagram of
 152 Figure 4(a)-(b). The parameters used for the power circuit are the following:

¹The inductor current never drops to zero.

153 $\delta_C = \delta_L = 0.1$ such that the switching period T is small compared to the
 154 circuit time constants, $I_r = 0.6$ is arbitrarily chosen and $V_r = \frac{1}{2}$ since it is
 155 required that $v_C = \frac{V_{in}}{2}$. The curves in red represent the current values at
 156 which the duty cycle saturates. Since at steady state we have $x_v[n] = V_r$ for
 157 all k_i , then (16) and (17) imply that $d_1[n] = d_2[n]$ and hence both duty cycles
 158 have similar upper and lower saturation limits, hereafter called boundaries
 159 (represented in red). These limits occur when $d_k[n] = 1$ and when $d_k[n] = 0$
 160 that is

$$d_k[n] = 1 \Leftrightarrow x_{i,sat,upper} = \frac{1}{k_i} + I_r \quad (23)$$

$$d_k[n] = 0 \Leftrightarrow x_{i,sat,lower} = I_r \quad (24)$$

161 2.2.2. Simulating the complete model

162 The fixed point of the complete accurate model is obtained from the equa-
 163 tion: $x^* = \Phi(\mathbf{d}[n])x^* + \Psi(\mathbf{d}[n])$ which implies that: $x^* = [(\mathbb{I} - \Phi(\mathbf{d}^*))^{-1}\Psi(\mathbf{d}^*)$
 164 provided that $[(\mathbb{I} - \Phi(\mathbf{d}^*))^{-1}]$ is a nonsingular matrix. Therefore, in the case of
 165 the complete model, we have no hope in solving this set of equations analyt-
 166 ically to obtain the fixed point x^* because \mathbf{d}^* is imposed by the controller to
 167 depend on x^* as $\mathbf{d}^* = K(x^* - X_r)$. Therefore, we will only present numerical
 168 simulations with similar controller parameters so comparison can be held.

169 When the proportional gains are fixed to obtain normal operating behav-
 170 ior (period-one behavior), numerical simulations have shown that the simple
 171 model and the accurate model have very similar responses (compare Fig-
 172 ure 4(a) with Figure 4(c) and Figure 4(b) with Figure 4(d)). According to
 173 the bifurcation diagrams, the bifurcation of the solution (from period-one
 174 to period-two) occurs when $k_i = 19$ for the simple model, however with the

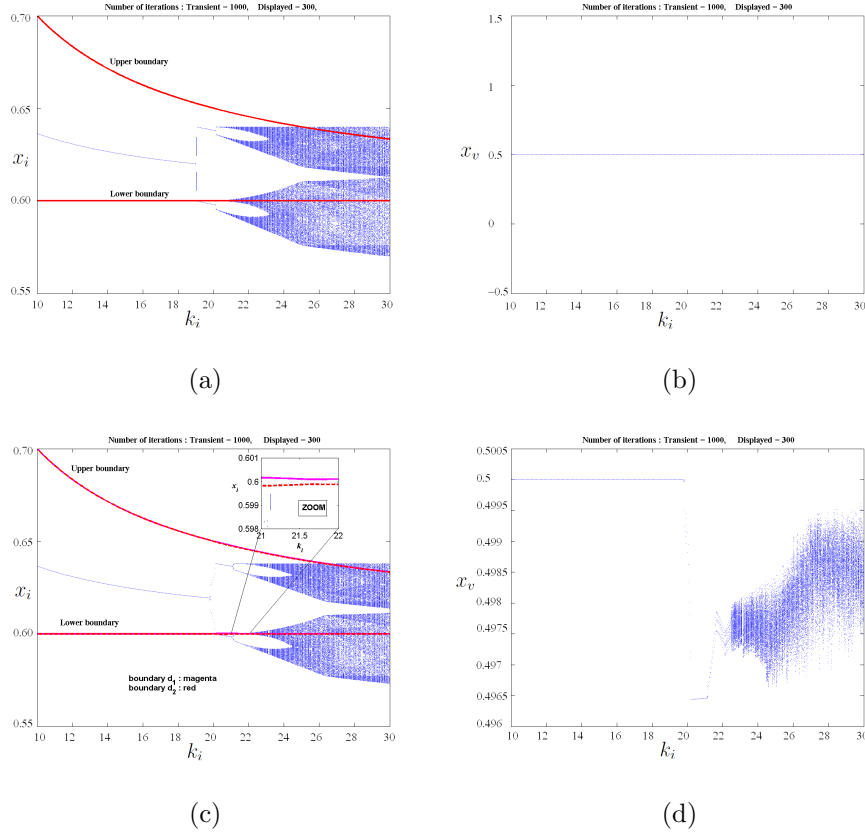


Figure 4: Bifurcation diagrams of x_i and x_v of the exact and simplified models with respect to bifurcation parameter k_i . (a) Inductor current of the simplified model. (b) Capacitor voltage of the simplified model. (c) Inductor current of the exact model. (d) Capacitor voltage of the exact model.

175 accurate model, the bifurcation occurs with larger value of $k_i \approx 19.8$. It is
 176 also worth to note that beyond the bifurcation value of k_i , there is no exact
 177 convergence of v_C to V_r (see Figure 4(d)), *i.e.*, we have also a chaotic behav-
 178 ior of the voltage which is practically more reasonable than having chaotic
 179 behavior of the current and stable behavior of the voltage as it has been
 180 obtained by the simple model. Consequently, calculating the boundaries in

181 terms of k_i is not trivial and the boundary expressions depend on the error
 182 $(x_v - V_r)$. To hold the comparison over, we have sketched the boundaries
 183 with the average value of the error. Although, those boundaries have no per-
 184 tinance, they may reveal explanation of some bifurcations as it can be seen
 185 in the zoomed part of Figure 4(c).

186 The comparison between the simulation results of the simple and the
 187 accurate models did not reveal a significant difference in the response. As
 188 a conclusion, we can consider the simple model for theoretical analysis and
 189 design, all along caring to check the performances of the suggested control
 190 method on the accurate model.

191 **3. Time-delayed feedback control (TDFC)**

192 Since we have noticed that, for the simple model, only the current be-
 193 havior undergoes two bifurcations until it reaches chaos, then it is tenable to
 194 apply the TDFC part in terms of current only. Omitting the saturation, the
 195 duty cycle expressions with TDFC are

$$d_1[n] = k_i(x_i[n] - I_r) + k_v(x_v[n] - V_r) + \eta_i(x_i[n] - x_i[n-1]) \quad (25)$$

$$d_2[n] = k_i(x_i[n] - I_r) - k_v(x_v[n] - V_r) + \eta_i(x_i[n] - x_i[n-1]) \quad (26)$$

The application of the TDFC controller increases by one the order of the discrete dynamical system. It is worthy to note here that this is one great advantage of using a discrete model. When applying TDFC to a continuous-time dynamical system, the system becomes infinite dimensional. Let's assume

the following notations

$$x[n] = \begin{pmatrix} x_1[n] \\ x_2[n] \\ x_3[n] \end{pmatrix} = \begin{pmatrix} x_v[n] \\ x_i[n] \\ x_i[n-1] \end{pmatrix} \quad (27)$$

Then the closed loop system, obtained by applying (25) and (26) to (15), is described by the following recurrent map

$$\begin{aligned} x_1[n+1] &= x_1[n] - 2k_v\delta_C(x_1[n] - V_r)x_2[n] \\ x_2[n+1] &= 2k_v\delta_L\left(x_1[n] - V_r - \frac{1}{2}\right)x_1[n] + (1 - \delta_L(k_i + \eta_i + 1))x_2[n] \\ &\quad + \delta_L\eta_ix_3[n] + \delta_L(1 + k_iI_r + k_vV_r) \\ x_3[n+1] &= x_2[n] \end{aligned} \quad (28)$$

The fixed point x^* of the normal operating mode is obtained by equating $x^*[n+1] = x^*[n]$. Since this equation gives $x_3^* = x_2^*$, then the TDFC part will vanish and we obtain similar fixed point as in the previous section, that is

$$(x_1^*, x_2^*, x_3^*) = \left(V_r, \frac{1 + k_iI_r}{1 + k_i}, \frac{1 + k_iI_r}{1 + k_i}\right) \quad (29)$$

196 3.1. Linear stability analysis

Linearizing system (28) around its fixed point yields to the following Jacobian matrix

$$J = \begin{pmatrix} 1 - 2k_v\delta_Cx_2 & -2k_v\delta_C(x_1 - V_r) & 0 \\ k_v\delta_L(2V_r - 1) & 1 - \delta_L(k_i + \eta_i + 1) & \delta_L\eta_i \\ 0 & 1 & 0 \end{pmatrix} \quad (30)$$

When substituting $x = x^*$ and $V_r = \frac{1}{2}$, the characteristic equation of the Jacobian matrix is

$$(\lambda - 1 + 2k_v\delta_Cx_2^*)(\lambda^2 - (1 - \delta_L(k_i + \eta_i + 1))\lambda - \delta_L\eta_i) = 0 \quad (31)$$

Clearly, $\lambda_1 = 1 - 2k_v\delta_C x_2^*$ is an eigenvalue for the system that we should place inside the unit circle of the complex plane by tuning the control gains k_i and k_v [16], [17]. Therefore, we have

$$|1 - 2k_v\delta_C x_2^*| < 1 \quad (32)$$

We finally may select k_v as follows

$$0 < k_v < \frac{1}{\delta_C x_2^*} \quad (33)$$

To obtain a fast response, we may impose the eigenvalue to be at the origin [16], [17]. This can be achieved by choosing

$$k_v = \frac{1}{2\delta_C x_2^*} \quad (34)$$

For the quadratic part of (31), we may apply Jury's criterion of stability (see [16]) that requires the following conditions to be satisfied

$$|a_0| < 1 \quad (35)$$

$$1 + a_1 + a_0 > 0 \quad (36)$$

$$1 - a_1 + a_0 > 0 \quad (37)$$

197 where

$$a_1 = -(1 - \delta_L(k_i + \eta_i + 1)) \quad (38)$$

$$a_0 = -\delta_L\eta_i \quad (39)$$

The above inequalities define several regions in the plane $k_i - \eta_i$. The region of stability is then the intersection between each defined region. Condition (35) gives

$$|\eta_i| < \frac{1}{\delta_L} \quad (40)$$

Condition (36) gives

$$k_i > -1 \quad (41)$$

which is satisfied when restricting the design to positive values of k_i and finally (37) yields

$$\eta_i < -\frac{k_i}{2} + \frac{1}{\delta_L} - \frac{1}{2} \quad (42)$$

198 The intersection of the regions defined by the inequalities (40)-(42) is the shaded area depicted in Figure 5. We notice that the critical proportional

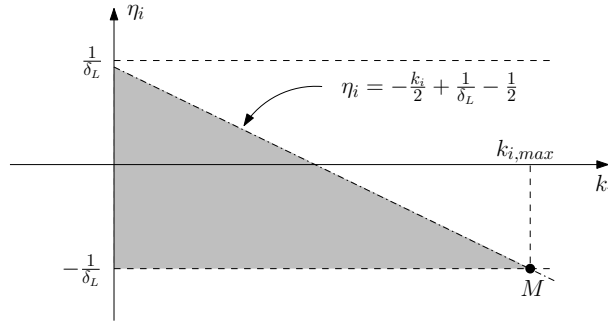


Figure 5: The stability zone in the parameter plane $k_i - \eta_i$.

199

200 gain k_i can be made larger, and then the stability region can be enlarged,
 201 by using a negative TDFC action ($\eta_i < 0$). However, we cannot increase the
 202 current gain k_i indefinitely. The intersection point M in Fig. 5 defines the
 203 maximum value of k_i which is given by $k_{i,max} = \frac{4}{\delta_L} - 1$.

204 3.2. Nonlinear stability analysis

Similar results can be obtained using a nonlinear stability analysis. We should note that the stability analysis we will present herein only concerns the non saturating case. A thorough analysis including the saturation of the duty cycles will be included in future work. To tackle our analysis, we define

the error vector as $e_1[n] = x_1[n] - V_r$, $e_2[n] = x_2[n] - I_r$ and $e_3[n] = x_3[n] - I_r$ and thus we obtain the error system

$$\begin{aligned} e_1[n+1] &= (1 - 2k_v\delta_C x_2[n])e_1[n] \\ e_2[n+1] &= 2k_v\delta_L\left(e_1[n] - V_r - \frac{1}{2}\right)e_1[n] + (1 - \delta_L(k_i + \eta_i + 1))e_2[n] \\ &\quad + \delta_L\eta_i e_3[n] + \delta_L(1 - I_r) \\ e_3[n+1] &= e_2[n] \end{aligned} \quad (43)$$

205 we remark that the error system is a nonlinear system in a cascade triangular
206 form

$$z_1[n+1] = f_1(z_1[n]) \quad (44)$$

$$z_2[n+1] = f_2(z_1[n], z_2[n]) \quad (45)$$

207 with $z_1 = e_1$ and $z_2 = [e_2, e_3]^t$. It has been proved that under mild condi-
208 tions, the stability of system (44)-(45) can be deduced from the simultaneous
209 asymptotic stability of the following subsystems [29]

$$z_1[n+1] = f_1(z_1[n]) \quad (46)$$

$$z_2[n+1] = f_2(z_1^*, z_2[n]) \quad (47)$$

where z_1^* is the fixed point of (46). Now, the first subsystem is described by

$$e_1[n+1] = (1 - 2k_v\delta_C x_2[n])e_1[n] \quad (48)$$

210 Hence, we may consider the Lyapunov function candidate $V(e_1[n]) = e_1[n]^2$.
211 In discrete systems we need to show that the first difference of the Lyapunov
212 function is negative *i.e.* $\Delta V(e_1[n]) = V(e_1[n+1]) - V(e_1[n]) < 0$.

$$\Delta V(e_1[n]) = 2k_v\delta_C x_2[n](2k_v\delta_C x_2[n] - 2)e_1[n]^2 \quad (49)$$

Knowing that $x_2[n]$ is always positive, then we can guarantee that ΔV is negative if the voltage gain k_v is chosen to satisfy condition (50).

$$0 < k_v < \frac{1}{\delta_C x_2[n]} = k_{v,cri} \quad (50)$$

213 We notice that condition (50) is equivalent to condition (33) around the fixed
 214 point. We can also optimize the choice of k_v to obtain fast convergence of
 215 the error to zero. This is done by minimizing $\Delta V(e_1[n])$, therefore we need
 216 to minimize $h(k_v) = 2k_v\delta_C x_2[n](2k_v\delta_C x_2[n] - 2)$ with respect to k_v and we
 217 get

$$k_{v,opt} = \frac{1}{2\delta_C x_2[n]} \quad (51)$$

218 As it can be observed, this condition is the same than that given in (34)
 219 with the difference that (34) is obtained from the linear approach with $x_2[n]$
 220 evaluated at the fixed point component x_2^* .

221 Until now, we can guarantee the stability of subsystem (48) at the origin
 222 by the choice of k_v , that is we obtain $e_1[n] = 0$ equivalent to $x_1[n] = V_r$. Now,
 223 by substituting $e_1[n] = 0$ in the second subsystem *i.e.* $z_2[n+1] = f_2(0, z_2[n])$
 224 we get

$$e_2[n+1] = (1 - \delta_L(k_i + \eta_i + 1))e_2[n] + \delta_L\eta_i e_3[n] + \delta_L(1 - I_r) \quad (52)$$

$$e_3[n+1] = e_2[n] \quad (53)$$

This is a linear system with characteristic equation

$$\lambda^2 - \left(1 - \delta_L(k_i + \eta_i + 1)\right)\lambda - \delta_L\eta_i = 0 \quad (54)$$

which has been studied in the previous section. To summarize, the analysis carried out here, we can say that if we choose $k_v < \frac{1}{\delta_C}$ and any pair of (k_i, η_i)

in the shaded area of Fig. 5, then the system evolves towards the fixed point (x_1^*, x_2^*, x_3^*) . To obtain a fast evolution, we might prefer to choose $k_v = \frac{1}{2\delta_C}$ and choose the pair of (k_i, η_i) that gives a minimum spectral radius that corresponds to eigenvalues of (54) with the least norm, that should be the pair that leads to equal real eigenvalues. The optimal value of η_i verifies:

$$\eta_{i,opt} = -\frac{1 + \delta_L(1 + k_i) - 2\sqrt{\delta_L(1 + k_i)}}{\delta_L} \quad (55)$$

225 3.3. Simulation Results of the TDFC Performances

226 In this section we present numerical simulations that confirm the theo-
 227 retical results obtained in the foregoing section. Following the theoretical
 228 analysis presented earlier, we have shown that k_v can be chosen to achieve
 229 $x_1 = V_r$, thus we have fixed $k_v = \frac{1}{2\delta_C} = 5$. Next, we will present simulations
 230 with different pairs of (k_i, η_i) .

231 First, in Fig. 6 we present the evolution of the current $x_2[n]$ and the
 232 voltage $x_1[n]$ with k_i increased near to its maximum $k_{i,cri} = 38.5$ and $\eta_i =$
 233 -9.85 thus the obtained steady state error is about 1%. To investigate the
 234 optimality relation presented in (55) we present two different curves of the
 235 current with $k_i = 31$ while η_i is fixed first at $\eta_i = -8$ and then at $\eta_i = \eta_{i,opt}$.
 236 The initial conditions were fixed at +5% of the fixed point values. The
 237 results are shown in Fig. 7. Clearly, the current settles to the steady state
 238 faster when using $\eta_{i,opt}$. To confirm the optimality relation, we present in
 239 Fig. 8 another simulation at $k_i = 25$.

240 Figure 9 depicts a 2D-bifurcation diagram where the behavior of the con-
 241 verter is elucidated for each pair of parameter (k_i, η_i) . This diagram is ob-
 242 tained from an ad hoc routine written in Matlab. The colors indicate different

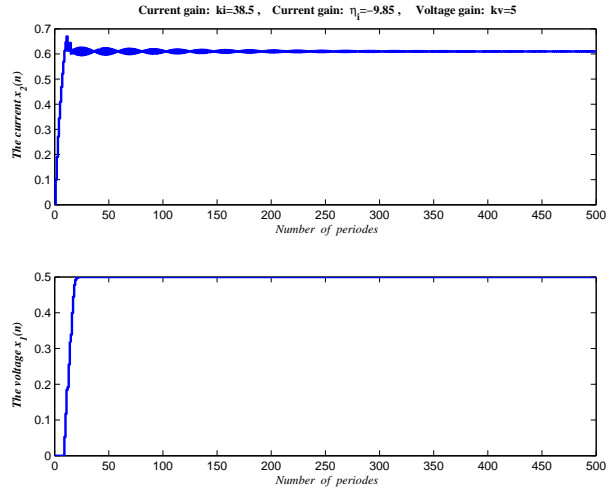


Figure 6: Evolution of the states under a Digital TDFC.

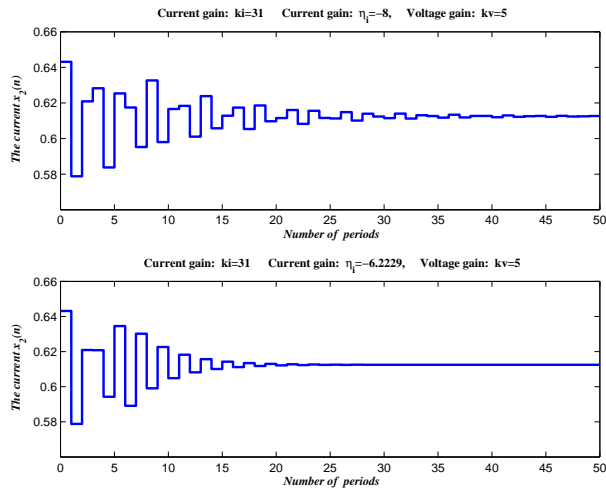


Figure 7: Evolution of the current under a Digital TDFC without and with optimized parameters.

243 periodic behaviors and black stands for chaotic motion. Brown stands for
 244 the T -periodic behavior and it confirms the triangular zone obtained theoret-
 245 ically. If we look at the horizontal line $\eta_i = 0$, we observe that the behavior

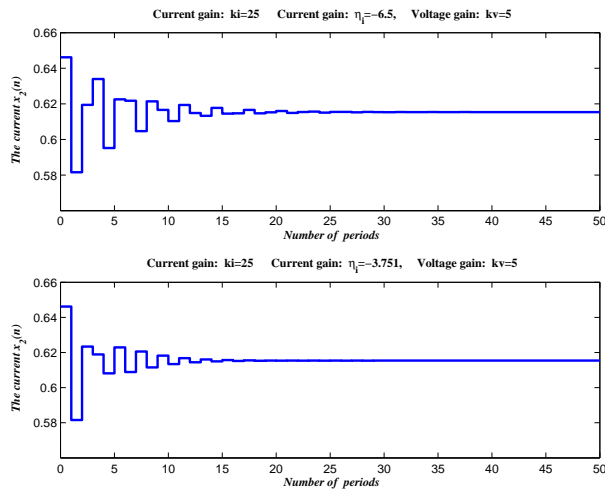


Figure 8: Evolution of the current under a Digital TDFC without and with optimized parameters (confirmation).

246 is exactly that obtained on Fig. 4 with red standing for $2T$ -periodic behavior
 247 and so on. In Fig. 9 we also plot in dashed line the curve of optimal pairs.
 248 Although, one may intuitively choose a pair situated in the middle of the
 249 stability zone (*e.g.* results shown in Figs. 7 and 8) thinking to have faster
 250 response when situated far from the stability limits; we notice that the op-
 251 timality curve is in fact and surprisingly very close to the upper limit of the
 252 stability zone.

253 3.4. Unpredicted Behavior: Coexistence of Attractors

254 When exploring the optimality condition obtained in (55), we have chosen
 255 to start the simulation in the vicinity of the fixed point with the aim to ob-
 256 serve the effect of the parameter choice and eliminate any other effects such
 257 as transients due to far initial conditions. It is worth noting that we have
 258 obtained a rather unpredictable behavior for some very particular pairs. We

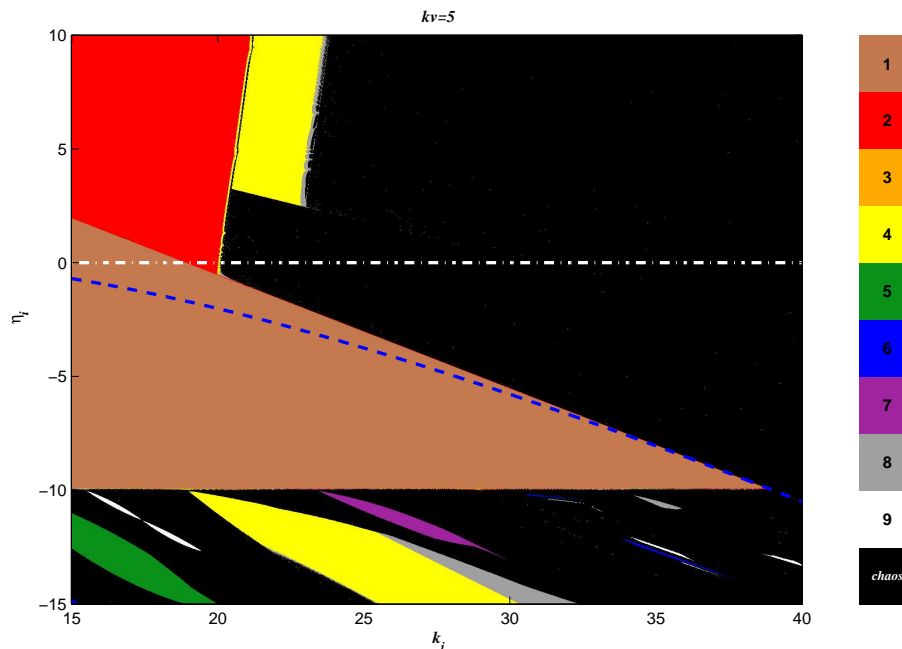


Figure 9: 2D-Bifurcation diagram of the converter behavior under the action of a Digital TDFC. Initial conditions simulations are selected at +5% of the fixed point values.

259 present in Fig. 10 the evolution of the current with $(k_i, \eta_i) = (31, -9)$. We no-
 260 tice that the current evolution is chaotic with values wandering within three
 261 different intervals (three piece chaotic attractor). Unlike the 2D-bifurcation
 262 diagram presented in Fig. 9 where simulations start at the origin of the state
 263 space, the 2D-bifurcation diagram presented in Fig. 11 is obtained by carry-
 264 ing simulations that start at +5% of the fixed point values. There, we have a
 265 slight difference in the region of $k_i = 31$. Indeed, chaotic behavior is obtained
 266 from pairs that according to the stability analysis, should give stable fixed
 267 point! We should point out here that this specific region changes slightly its
 268 shape according to the initial conditions and the value of k_v . To make an in
 269 depth investigation we plotted a bifurcation diagram with $\eta_i = -9.8$ using

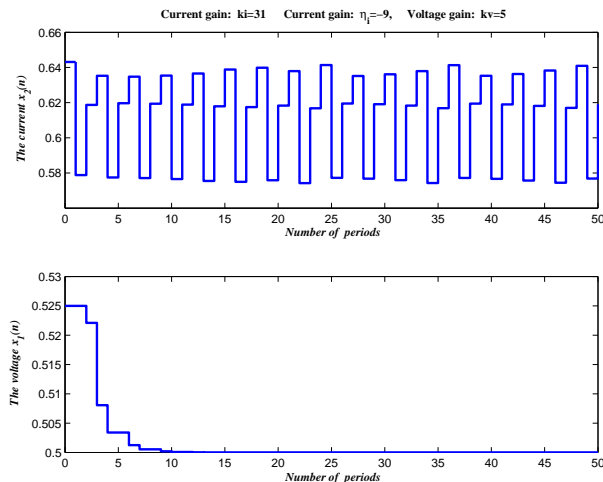


Figure 10: Evolution of the current $x_1[n]$ and the voltage $x_2[n]$ under a Digital TDFC with particular parameters. Initial conditions simulations are selected at the fixed point values.

270 simulations that start at $+5\%$ of the fixed point values. The diagram of
 271 Fig. 12 shows a chaotic attractor with three separate regions in the interval
 272 $(29.3, 33.5)$ of the bifurcation parameter k_i . The sensitivity of the current
 273 evolution to initial condition implies here the coexistence of periodic motion
 274 and chaotic behavior. In Fig. 13, we present the basins of attraction of the
 275 fixed point (in brown) and of the chaotic attractor (in black) using some
 276 fixed parameters. We first note that the basins of attraction are intertwined.
 277 The separatrices are mainly vertical; making the basins independent of $x_3[n]$,
 278 except in some region where they have a different shape. Actually, in the
 279 regions where the basins of attraction are formed from vertical strips, we
 280 note that the duty cycles are both saturated either at zero $d_1 = d_2 = 0$ for
 281 lower values of initial conditions (call that region \mathcal{R}_0) or saturated at one
 282 $d_1 = d_2 = 1$ for higher values of initial conditions (call that region \mathcal{R}_1). The

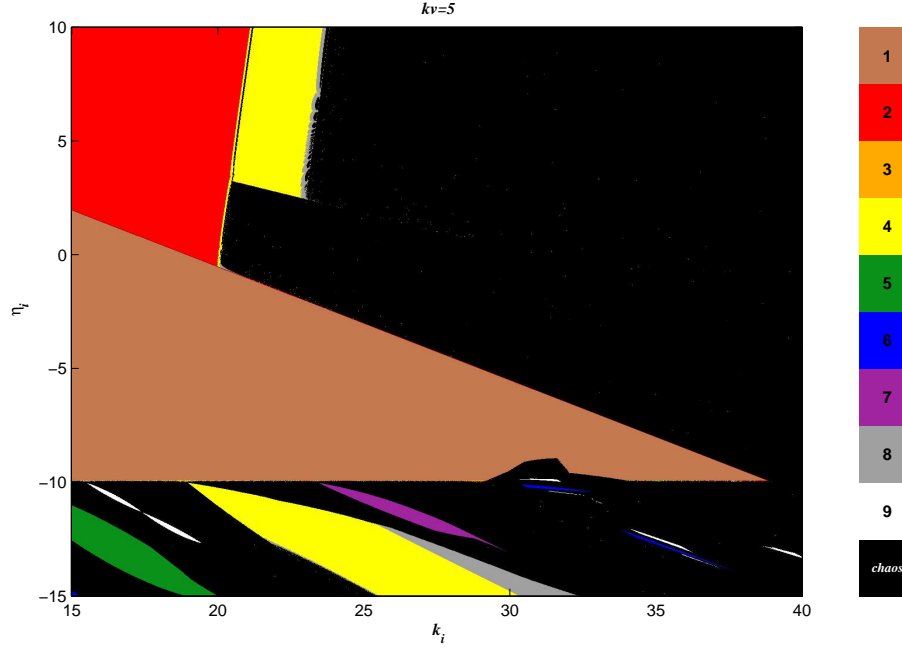


Figure 11: 2D-Bifurcation diagram of the converter behavior under the action of a Digital TDFC. Initial conditions simulations are selected at +5% of the fixed point values. White dashed line corresponds to the bifurcation 1D-diagram of Fig. 4(a). Blue dashed line corresponds to the optimum value of η_i in terms of k_i .

283 boundaries of saturation of d_1 and d_2 are plotted in yellow and red respec-
 284 tively.

285 Now, when the system starts evolving in \mathcal{R}_0 , it is attracted towards a
 286 virtual fixed point in \mathcal{R}_1 and vice-versa. The evolution of the current follows
 287 the line $\mathcal{L}_0 : x_3[n] = \frac{1}{1-\delta_L}x_2[n] - \frac{\delta_L}{1-\delta_L}$ in \mathcal{R}_0 (see the yellow circles and white
 288 asterisks). In \mathcal{R}_1 the evolution of the current follows the line $\mathcal{L}_1 : x_3[n] =$
 289 $\frac{1}{1-\delta_L}x_2[n]$ (see the cyan triangles and magenta diamonds). Eventually, the
 290 evolution should reach a region delimited by the duty cycle saturation lines
 291 and the lines \mathcal{L}_0 and \mathcal{L}_1 (call that region \mathcal{R}_{fp}). Once the trajectories are

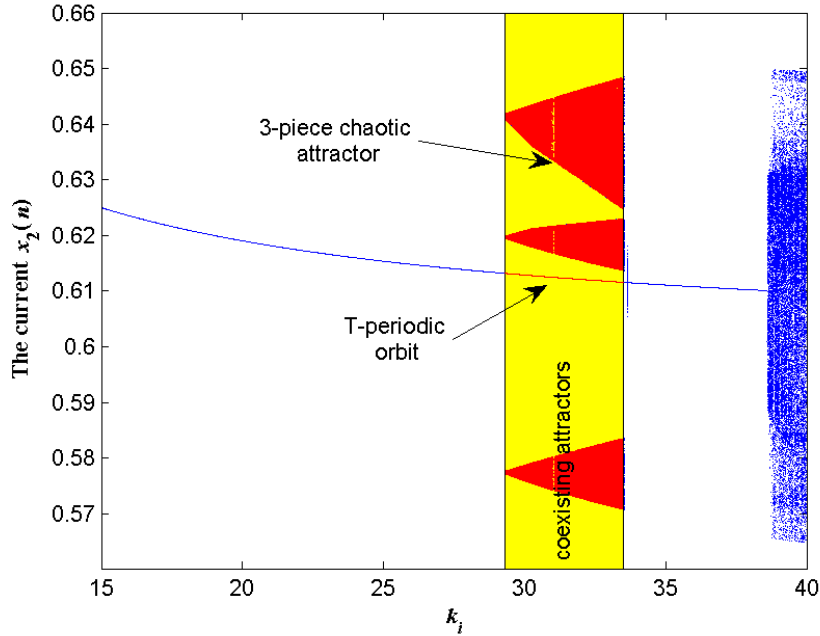


Figure 12: Bifurcation diagram showing coexistence of chaotic attractor and periodic behavior under the action of a Digital TDFC. Three piece chaotic attractor coexists with T -periodic orbit within the range $(29.3, 33.5)$ of the bifurcation parameter k_i .

292 in \mathcal{R}_{fp} , the system will either converge to the fixed point or to the three-
 293 regions chaotic attractor. In Fig. 14, we present a zoom of the region \mathcal{R}_{fp} ,
 294 there we observe that trajectories attracted by the fixed point will reach it
 295 while spiraling around it since we can check that the eigenvalues associated
 296 to the fixed point are complex conjugate. Now, in absence of saturation of
 297 the duty cycles, the fixed point is globally attractive. However, in presence
 298 of the saturation, points which are still far from the fixed point, *i.e.* next to
 299 the boundaries of the region \mathcal{R}_{fp} , might be attracted to another virtual fixed
 300 point following different trajectories. This is the key point of the problem. If

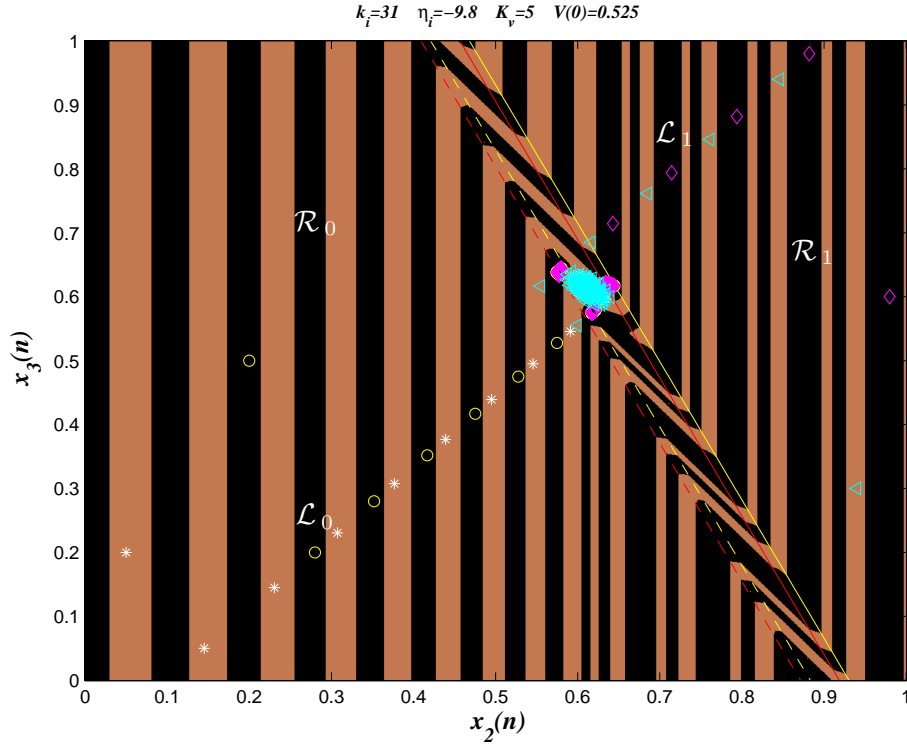


Figure 13: Riddled basins of attraction of the coexisting attractor and the fixed point. In brown, the basin of attraction of the fixed point and in black the basin of attraction of the there-regions chaotic attractor. In white asterisks and cyan triangles we present two system trajectories converging to the fixed point and initiated at \mathcal{R}_0 and \mathcal{R}_1 respectively. In yellow circles and magenta diamonds we present two system trajectories converging to the chaotic attractor and initiated at \mathcal{R}_0 and \mathcal{R}_1 respectively.

301 a trajectory gets close to the fixed point such that its future evolution does
 302 not leave the non-saturation region then it converges to the fixed point. On
 303 the other hand, if a trajectory approaches the fixed point but still relatively
 304 far so that its future steps are outside the non-saturation region then it will
 305 go farther by the effect of the virtual fixed points but cannot leave the region
 306 \mathcal{R}_{fp} . That's how a new chaotic attractor has been created.

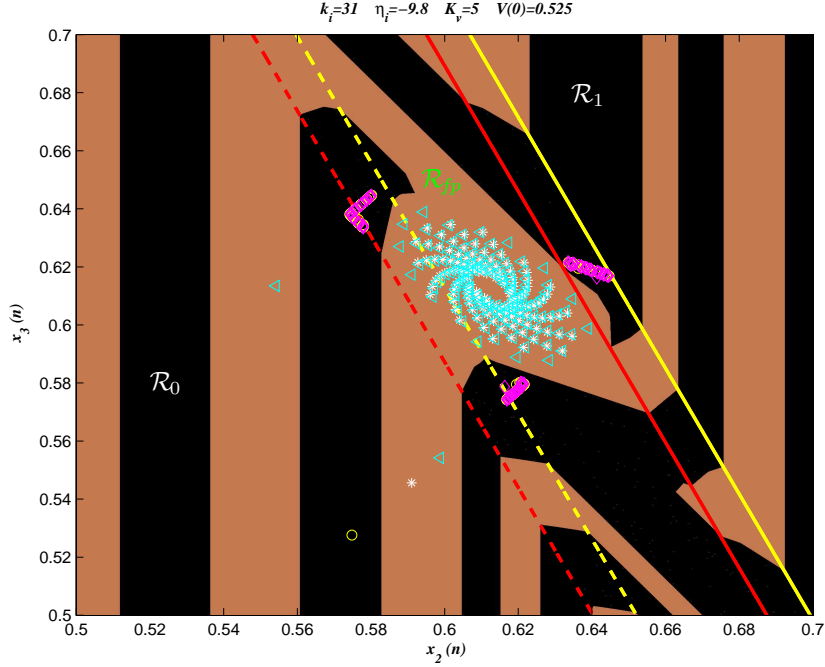


Figure 14: Zoom about the fixed point and the coexisting chaotic attractor. In brown, the basin of attraction of the fixed point and in black the basin of attraction of the three-regions chaotic attractor. In white asterisks and cyan triangles we present two system trajectories converging to the fixed point and initiated at \mathcal{R}_0 and \mathcal{R}_1 respectively. In yellow circles and magenta diamonds we present two system trajectories converging to the chaotic attractor and initiated at \mathcal{R}_0 and \mathcal{R}_1 respectively.

307 Conclusion

308 In this paper, we have presented two discrete models of a two-cell DC-
 309 DC buck converter, starting from the most complete and complicated one
 310 to the approximated and simplest one. The advantage of the complex one
 311 is to reflect with a great fairness the behavior of the real system, though it
 312 cannot be processed from an analytical point of view. On the contrary, the

313 advantage of the simple and approximated model is that it can be handled
314 analytically with the aim to design controller to harness any abnormal be-
315 havior that can be exhibited by the converter. The obtained models have
316 been then simulated numerically to predict behaviors of the converter. The
317 simplicity of its architecture does not prevent it from having a rich spectrum
318 of complex dynamical behavior. The most amazing among these behaviors is
319 the chaotic one which is usually classified as undesirable in power electronics.
320 The purpose of the present work has been to design a feedback control signal
321 that help to suppress the chaotic behavior with the aim to ameliorate the
322 performance of the converter and to enlarge the domain of its application.
323 Here, we have presented the TDFC controller. The normal periodic behavior
324 has been stabilized in some parameter space where chaotic behavior has been
325 observed. The controller presents several advantages. The first is that it still
326 has a simple expression to be implemented physically. The second advantage
327 is the speed of its action if the time delayed feedback gain η_i is optimized
328 and the third advantage is its non-invasiveness since the TDFC term van-
329 ishes when the system reaches the period-one behavior. Of course, we do
330 not claim to present the perfect solution to the problem in question, albeit
331 several advantages have been cited. In fact, from a control theory point of
332 view, our analysis has neglected the duty cycle saturation which requires an
333 in depth analysis to guarantee the stability and to avoid any bad surprise in
334 the real operation of the converter like the coexistence of attractors discov-
335 ered by numerical simulations. This constitutes a first point that we have
336 already started to investigate in order to present a complete theoretical anal-
337 ysis about this issue. This complete study together with an experimental

338 validation will be the subject of a future work.

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